

Starlight — time and again

Samuel Conner & Don Page

Dr Humphreys rejects our Newtonian illustration and relativistic demonstration^{1,2} of the identical gravitational fields of bounded and unbounded universes on the basis of his claim that unbounded universes cannot possess the property of spherical symmetry.^{3,4} Humphreys' claims are mistaken, and are based on a misunderstanding of the physical and mathematical meaning of spherical symmetry. An object is said to possess the property of spherical symmetry with respect to a particular point if it is unchanged by an arbitrary rotation about an arbitrary axis passing through that point. For objects which have a spherical boundary, there is only one such point of spherical symmetry,⁵ and this point constitutes a unique centre to the object. Humphreys believes that objects which do not possess a spherical boundary, such as the unbounded universe models of 'big bang' cosmology, cannot possess spherical symmetry. However, such unbounded universes in fact are spherically symmetric about *every* point: an arbitrary rotation about an arbitrary axis through any point will leave the system unchanged. The spherical symmetry of the unbounded 'big bang' models is also obvious from the manifest spherical symmetry of the Robertson-Walker form of the metric,⁶ which is commonly employed to describe their geometry.

Spherical symmetry with respect to a particular point is identical to the notion of *isotropy* about that point. Unbounded standard 'big bang' universes are isotropic with respect to every point in their interiors, which is to say that they are spherically symmetric from the point of view of every point. Such universes are not acentric, as claimed by Humphreys, but rather have infinitely many centres of spherical symmetry. The centres of spherical symmetry are not unique, but that does not mean that they do not exist at all. Humphreys' claim that the Copernican principle is incompatible with spherical symmetry⁷ is another misunderstanding. The Copernican principle imposes spherical symmetry with respect to every point because it imposes the property of isotropy at every point. The spherical symmetry of standard unbounded 'big bang' models is discussed, using the language of isotropy, in every cosmology textbook. Readers interested in a thorough discussion are referred to the chapter which Weinberg devotes to his discussion of symmetric spaces⁸ and especially to the conclusion of that chapter, which discusses the special case of 'Spherically Symmetric Homogeneous Spacetime' and derives the Robertson-Walker form of the metric solely

from the symmetry properties of this geometry.⁹

We did not commit a 'big blunder,' as Humphreys charges,¹⁰ in our appeal to the obvious symmetry properties of unbounded homogeneous and isotropic matter distributions. Our appeal to these properties is valid and the conclusions, we have drawn from these properties regarding the identity of the gravitational properties of Humphreys' bounded universe and the unbounded 'big bang' universe are correct.

Another demonstration of the identity of the gravitational properties of the bounded and unbounded universes is found in consideration of the manner in which each of these models decelerates as it expands. It is straightforward, using simple Newtonian arguments,¹¹ to show that the thin shell of matter at the boundary of Humphreys' bounded universe experiences a radial deceleration of

$$\frac{d^2 D_B(\tau)}{d\tau^2} = -D_B \frac{4}{3} \pi G \rho \quad (1)$$

Where ρ is the mass density of the bounded universe¹² at time τ , $D_B(\tau)$ is the distance from the centre of the matter sphere to the expanding boundary, and t is the time measured by an observer at rest at the centre of the matter sphere.¹³ This calculation may also be employed to calculate the deceleration of any matter shell located at the physical distance D from the center: just substitute D for D_B in equation 1.

To compare the deceleration of the unbounded universe to that of the bounded universe, we will consider the deceleration of a comoving shell of matter located at the same physical distance from the adopted origin of coordinates as the boundary (or any other choice of matter shell) of the bounded matter sphere is from its unique centre. Denote the comoving radial coordinate of this matter shell by η_{shell} and the physical distance to the origin by $D_{shell, unbounded}(\tau)$. The metric tells us how to compute $D_{shell, unbounded}(\tau)$:

$$D_{shell, unbounded}(\tau) = a(\tau) \int_{\eta=0}^{\eta=\eta_{shell}} \frac{d\eta}{\sqrt{1-k\eta^2}} \quad (2)$$

The deceleration of this shell is simply

$$\frac{d^2 D_{shell, unbounded}(\tau)}{d\tau^2} = \frac{d^2 a(\tau)}{d\tau^2} \int_{\eta=0}^{\eta=\eta_{shell}} \frac{d\eta}{\sqrt{1-k\eta^2}} = D_{shell, unbounded}(\tau) \frac{1}{a(\tau)} \frac{d^2 a(\tau)}{d\tau^2} \quad (3)$$

The deceleration of $a(\tau)$ is given by the Friedman deceleration equation^{14,15} which, for the cosmic matter content under consideration, is simply

$$\frac{1}{a(\tau)} \frac{d^2 a(\tau)}{d\tau^2} = -\frac{4}{3} \pi G \rho \quad (4)$$

This relation may be combined with equation 3 to give

$$\frac{d^2 D_{shell,unbounded}(\tau)}{d\tau^2} = -D_{shell,unbounded}(\tau) \frac{4}{3} \pi G \rho \quad (5)$$

This equation is obviously identical to the bounded case deceleration equation (equation 1). Matter shells located at the same distance from the adopted origin of coordinates experience the same deceleration, regardless of whether they are in a bounded or unbounded universe. This deceleration is manifestly gravitational in origin, since there are no forces other than gravity acting on the shells.¹⁶ In Newtonian language, gravitational accelerations are due to the ‘gravitational field.’ Therefore, the gravitational fields in the interiors of the bounded and unbounded universes are identical.¹⁷ This simple illustration overthrows the central claim of Humphreys’ cosmological theorizing.

To understand the cause of the sign change in the Klein form of the metric, it is helpful to understand where the Klein metric components come from. Unlike the Robertson-Walker form of the metric, the Klein form of the metric is not derived from first principles using symmetry properties or the field equations of general relativity applied to a bounded matter distribution. Rather, the vacuum Schwarzschild coordinate system is extended inward from the surface of the matter and the resulting coordinate system (including the imaginary part of t_{Klein}) is used to transform the Robertson-Walker metric components¹⁸ into the Klein coordinate system.¹⁹ The transformation is given by the conventional tensor transformation relation.²⁰ Making the transformation relation explicit, one has

$$\beta(t,r) = g_{tt,Klein} = (g_{Klein}^t)^{-1} = \left(g_{RW}^{\tau\tau} \left(\frac{\partial t_{Klein}}{\partial \tau} \right)^2 + g_{RW}^{\eta\eta} \left(\frac{c \partial t_{Klein}}{\partial \eta} \right)^2 \right)^{-1} \quad (6)$$

It is a straightforward mathematical exercise to show that this equation, using the full complex form of $t_{Klein}(\tau, \eta)$, gives Klein’s metric component $\beta(t,r)$. It is also easy to show that Equation 6 simplifies to

$$\beta(t,r) = \left(\frac{\partial t_{Klein}}{\partial \tau} \right)^{-2} \left(\frac{1-\eta^2}{1-\frac{a_{max}}{a} \eta^2} \right) = \left(\frac{\partial t_{Klein}}{\partial \tau} \right)^{-2} (1-\eta^2) a(t,r) \quad (7)$$

where $a(t,r)$ is the Klein g_{rr} metric component. It is clear from equation 7 that $\beta(t,r)$ switches sign whenever $a(t,r)$ does. These sign changes are uninteresting, since the metric component signs have a Lorentzian appearance on both sides of such a sign change surface. The change which Humphreys considers interesting is when $\beta(t,r)$ switches sign but $\alpha(t,r)$ does not. It is clear from equation 7 that this type of change occurs if, and only if, $(\partial t_{Klein} / \partial \tau)^2$ changes sign;²¹ that is, if, and only if, $(\partial t_{Klein} / \partial \tau)$ changes from real to imaginary or vice versa.

A number of consequences inescapably follow from this fact. First, the ‘interesting’ sign change in β is simply an artefact²² of the change of the Klein time coordinate

dt_{Klein} from real to imaginary. This is a trivial form of metric component sign change and it has no physical consequences.²³ Second, Humphreys’ claim that the integral from which $t_{Klein}(t, \eta)$ is computed ‘should only be evaluated for values of the variable which are real, not imaginary,’^{24,25} eliminates the sign change in $\beta(t,r)$, since the sign change is caused by the change from real dt_{Klein} to imaginary dt_{Klein} . Finally, we note that Humphreys’ alleged restriction on t_{Klein} is not present in the published research literature on the Klein form of the metric.²⁶

Finally, we offer the following observations on Humphreys’ appeals to the research literature on classical signature change. First, Humphreys fails to note the speculative character of this literature. No one knows at present whether, and, if so, under what physical conditions classical signature change may occur.²⁷ This literature certainly does not establish Humphreys’ claims (and, in any event, Humphreys’ model does not undergo the signature change described in this literature; the sign change in Humphreys’ model is the consequence of the imaginary character of the Klein time coordinate). Second, Humphreys misunderstands the criterion for signature change proposed by Ellis, *et al.*²⁸ This criterion relates to the local non-gravitational energy density of the universe. Ellis *et al.* propose that, if the details of this energy content are such that the dynamics of the universe would lead to da/dt imaginary, then signature change should take place to keep da/dt real.²⁹ Humphreys erroneously includes gravitational potential energy in the local energy budget,³⁰ when in fact the only contributions to da/dt are matter fields, spacetime curvature and the cosmological constant.³¹ Third, Humphreys fails to note that much of the published literature on signature change³² applies to unbounded as well as bounded matter distributions,³³ which shows that it is not necessary to posit a boundary for signature change to occur.

Finally, Humphreys erroneously claims (and uses this false claim as justification for his erroneous rejection of the Robertson-Walker form of the metric) that signature change of the type considered by Ellis *et al.* cannot take place in the spacetime described by the Robertson-Walker metric. In fact, signature change will take place in the Robertson-Walker metric if the Ellis, *et al.* criterion (assuming it to be valid) is satisfied.³⁴ This is easy to show by writing the Robertson-Walker metric with the a -dependence of the cosmic time coordinate t made explicit:

$$ds^2 = c^2 \left(\frac{d\tau}{da} \right)^2 da^2 - a^2 \left(\frac{d\eta^2}{1-k\eta^2} + \eta^2 (d\theta^2 + \sin^2 \theta d\phi^2) \right) \quad (8)$$

The radius of curvature, a , of the universe is by definition a real number, so da is necessarily real. The signature of this metric will change from Lorentzian to Euclidean if the expansion rate of the universe, da/dt , changes from real to imaginary. The actual behaviour of

the expansion rate is determined by the dynamics produced by the matter/energy content of the universe, an issue which is independent of the employment of the Robertson-Walker form of the metric (the R-W metric is valid for locally homogeneous and isotropic universes, regardless of the details of the expansion dynamics). For the Friedman-Lemaitre class of R-W cosmologies (pressureless dust with non-zero cosmological constant), which are an excellent description of the late-time behaviour of the real universe, the expansion rate is given by the Friedman equation³⁵ for the Hubble parameter:

$$\frac{da}{d\tau} = aH_0 \left[\Omega \left(\frac{a_0}{a} \right)^3 + (1 - \Omega - \Omega_\Lambda) \left(\frac{a_0}{a} \right)^2 + \Omega_\Lambda \right]^{\frac{1}{2}} \quad (9)$$

where a_0 is the present radius of curvature of the universe, H_0 is the present value of the Hubble parameter, about 75 km/s/Mpc, and Ω and Ω_Λ are, respectively, the present matter and vacuum energy densities, in units of the critical density. This expansion rate can, in principle, be imaginary in the past (that is, for $a < a_0$) if Ω_Λ is much larger than Ω ; in such a case, the Robertson-Walker metric can have Euclidean signature for a range of past a . This shows the falsity of Humphreys' claim that the Robertson-Walker metric is not general enough to include the possibility of signature change. In fact, contrary to Humphreys' claims, the Robertson-Walker metric is every bit as general as the modified version employed by Ellis, *et al.* Indeed, the two forms of the metric are transformed into each other by the simple coordinate transformation³⁶ $(cdt/da)^2 = N(t)$, $da = dt$. Just as the Klein form of the metric is simply a different coordinate representation of the same underlying geometry described by the Robertson-Walker form of the metric, so also is the modified metric employed by Ellis, *et al.* Humphreys' claim that the geometry described by Ellis, *et al.* is profoundly different from the geometry of the Robertson-Walker metric is mistaken: the two geometries are identical.³⁷

Having shown that the Robertson-Walker form of the metric does, in principle, permit classical signature change, the question remains, 'has the universe actually experienced this in the past?'

The answer to this question is almost certainly 'not in the observable history of the universe (that is, not since the cosmic microwave background radiation decoupled from the matter of the universe)'. The reason for this is that the present equation of state of the matter content of the universe repels the universe from any Euclidean region which may exist: as the scale factor a approaches the transition to Euclidean signature, the expansion or contraction slows to a stop and reverses itself. This is what causes negative energy 'big bang' models to stop expanding and positive energy models with too large a value of Ω_Λ to 'bounce' in the past. The Euclidean region is a classically forbidden region of, essentially, negative cosmic kinetic energy (recall that the expansion rate da/dt is imaginary,

so that $(da/dt)^2$ is negative). It is not known whether there may be other equations of state which would permit a transition from positive to negative $(da/dt)^2$. We know of no such proposals and the literature cited by Humphreys contains none.

It may be possible to invent unusual hypothetical equations of state which would allow a homogeneous and isotropic universe to undergo signature change. However, such unusual equations of state bear no resemblance to the actual equation of state of the known matter and energy content of the real universe. Further, if one were to adopt such an unusual description of the expansion dynamics and impose a hypothetical signature change surface at some point in the past, this still would not solve the light travel problem, for such a signature change would occur simultaneously (that is, at the same value of a and t) throughout the universe, so that there would be no differential ageing of different parts of the universe. This simultaneity is imposed by the fact that cosmic time is synchronous with the expansion in all locally homogeneous and isotropic universes, so that da/dt is the same function of a throughout the universe. Therefore, if da/dt changes from real to imaginary, this change will take place at the same value of a and the same cosmic time t throughout the universe. In addition, by reducing the proper time available for light propagation, such a scenario would reduce the distance to the particle horizon (the greatest distance light can travel since the beginning of the universe).

If the location of a hypothetical signature change surface were adjusted to provide only 6,000 years of proper time since the beginning of the universe, as proposed by Humphreys, the particle horizon would be only about 6,000 light years distant, making all objects more distant than this invisible to observers on Earth. In fact, the furthest visible objects have been measured to be on the order of 10 billion light years distant (measurements with which Humphreys concurs³⁸). This indicates that, regardless of the number and 'duration' of past episodes of metric signature change, at least 10 billion years of proper time have elapsed since the beginning of the universe. As we have noted previously, the observed properties of the universe and the validity of General Relativity as a description of its behaviour over past time are incompatible with a recent origin. Humphreys' appeal to signature change cannot solve the light travel problem.

References

1. Conner, S.R. and Page, D.N., Starlight and time is the big bang, *CEN Tech. J.* 12(2):177-179, 1998.
2. Conner, S.R. and Page, D.N., *The Big Bang Cosmology of Starlight and Time*, unpublished manuscript, ~ 200 pages, 1997. Readers interested in obtaining this document may contact Mr Conner by mail at 10 Elmwood Avenue, Vineland, N.J. 08360, USA or by email at CERS_corresp@hotmail.com.
3. Humphreys, D.R., New vistas of spacetime rebut the critics, *CEN Tech. J.* 12(2):203-205, 1998.

4. Humphreys, D.R., Letter to the Editor, *CEN Tech. J.* **13**(1):59, 1999.
5. There is only one point of global spherical symmetry for a bounded spherical object because there is only one point about which an arbitrary rotation will map the *surface* of the object into itself. If there is no surface, then it is possible for there to be many points about which the object has spherical symmetry.
6. The Robertson-Walker form of the metric may be written

$$ds^2 = c^2 d\tau^2 - a^2(\tau) \left[\frac{d\eta^2}{1 - k\eta^2} + \eta^2 (d\theta^2 + \sin^2 \theta d\phi^2) \right]$$

The spherical symmetry of this formula is evident in that the angular part of the metric is simply the customary spherical line element

$$d\theta^2 + \sin^2 \theta d\phi^2$$

7. Humphreys, Ref. 3, pp. 205–206.
8. Weinberg, S., *Gravitation and Cosmology*, John Wiley and Sons, chapter 13, ‘Symmetric Spaces,’ pp. 375–404, 1972.
9. Weinberg, Ref. 8, pp. 403–404.
10. Humphreys, Ref. 4, p. 59.
11. It should be noted that such Newtonian arguments are not strictly valid, but are only approximately valid in the limit of low particle velocities, weak gravitational fields and nearly flat spacetime geometry. The Newtonian result for cosmic deceleration, computed in this limit, matches the fully relativistic result derived from the Einstein field equations when the pressure contributions to the energy-momentum tensor are negligible compared with ρc^2 .
12. We confine our attention to the ‘cold dust’ universe models considered by Humphreys. The Newtonian analysis may be extended to more general cases of non-zero pressure and non-zero cosmological constant, but the resulting identity of the deceleration of bounded and unbounded universes is not affected, so the present simple case suffices to illustrate that identity.
13. In Newtonian gravitation, there is no such thing as ‘time dilation’ and one need not employ the metric to compute physical (that is, ‘proper’) distances. In the Newtonian limit of GR, the spacetime metric approaches the metric of flat spacetime, so that (for appropriately chosen coordinates) there is little difference between coordinate and physical distances or between coordinate and proper time intervals for a stationary observer. One may do a fully relativistic analysis of the deceleration of a bounded matter sphere; the result is identical to that for the unbounded matter distribution, which is described below.
14. Peebles, P.J.E., *Principles of Physical Cosmology*, Princeton University Press, p. 312, 1993.
15. Weinberg, Ref. 8, p. 472.
16. There is no pressure in ‘cold dust’ models.
17. This point has been made previously, in Ref. 2 and in Conner, S.R. and Ross, H.N., *The Unraveling of Starlight and Time*, 1999, available from Reasons to Believe, <www.reasons.org>.
18. which, *contra* Humphreys, are indeed a valid metric description of the geometry of the bounded matter sphere (see Weinberg, Ref. 8, pp. 342–345.). If the Robertson-Walker metric were not valid as applied to the bounded matter sphere, then the Klein/Weinberg method of deriving the Klein metric, by transforming the Robertson-Walker metric into the Klein coordinate system, would fail and the resulting Klein metric would itself not be a valid description of the bounded matter sphere geometry.
19. Weinberg, Ref. 8, pp. 345–346.

$$g_{\alpha\beta, Klein} = g_{\mu\nu, RW} \frac{\partial x_{\alpha, Klein}}{\partial x_{\mu, RW}} \frac{\partial x_{\beta, Klein}}{\partial x_{\nu, RW}}$$

The simplest implementation of this relation, and that chosen in this letter, uses the contravariant coordinate transformation $\frac{\partial x_{Klein}^\alpha}{\partial x_{RW}^\mu}$ and the contravariant metric components $g^{mm} = (g_{\mu\mu})^{-1}$:

$$g_{\alpha\alpha, Klein} = (g_{Klein}^{\alpha\alpha})^{-1} = \left((g_{\mu\mu, RW})^{-1} \left(\frac{\partial x_{Klein}^\alpha}{\partial x_{RW}^\mu} \right)^2 \right)^{-1}$$

21. The factor $(1 - \eta)^2$ is always greater than or equal to zero.
22. as noted in Conner, S.R., Letter to the Editor, *CEN Tech. J.*, **13**(1):56–58, 1999.
23. It is obvious that one can always switch the sign of g_{tt} simply by changing the time coordinate from dt to $\sqrt{-1} dt$, a coordinate transformation called ‘Wick rotation’. Such a change has no physical significance.
24. Humphreys, Ref. 3, p. 209.
25. Humphreys, Ref. 4, p. 59.
26. Weinberg’s discussion of the derivation of the Klein metric (Weinberg, Ref. 8, pp. 345–346.) is the most accessible English-language treatment. Weinberg’s discussion makes it clear that imaginary dt_{Klein} is an inescapable consequence of the extension of the Schwarzschild coordinate system into the interior of the matter sphere. Humphreys’ proposed restriction on Weinberg’s integral expression for t_{Klein} is unjustified, for it destroys the resulting form of $\beta(r, \tau)$. To obtain Klein’s form of $\beta(r, \tau)$, with its sign-switching behaviour, it is necessary to perform the full integral for t_{Klein} with no restrictions.
27. An example of an unresolved problem in this research is the open question of what type of boundary and continuity conditions should be imposed across a signature change surface and the fact that the Einstein field equations are not obeyed on the surface. Another open question is the meaning of the discontinuities in physical properties (such as particle rest masses) which occur across a signature change surface. It should be noted that there is no such discontinuity in particle masses or the cosmic mass/energy density in the Klein metric, which shows that the coordinate-artefact-induced metric component sign change which occurs in the Klein metric is unrelated to the hypothetical processes considered in the literature cited by Humphreys.
28. Ellis, G., Sumeruk, A., Coule, D. and Hellaby, C, Change of signature in classical relativity, *Classical and Quantum Gravity*, **9**:1546–1548, 1992.
29. Ellis, *et al.*, Ref. 28, p. 1548.
30. Humphreys, Ref. 3, p. 202. Humphreys claims that Ellis, *et al.*’s proposed signature change criterion is ‘closely related to the gravitational potential of our spherical, bounded, distribution of matter’, but in fact it is unrelated. The local character of the Ellis, *et al.* criterion is obvious from the fact that their signature change occurs at the same point in the cosmic expansion (at the same value of the cosmic scale parameter a), regardless of location in the universe, as expected for a locally homogeneous geometry, whether bounded or unbounded.
31. See, for example, the discussion in Peebles, Ref. 14, pp. 259–268, 280–313.
32. including the article by Ellis, *et al.*, Ref. 28, which Humphreys cites extensively.
33. In particular, Ellis *et al.* study signature change in the context of conventional unbounded cosmology.
34. Humphreys is mistaken in claiming that one must explicitly incorporate a sign switch in the g_{00} metric component in order for signature change to be possible. Transformation from Lorentzian to Euclidean signature

is a physical process which, if it is possible at all, will occur regardless of the coordinate representation used for the metric. Ellis, *et al.* make the signature transition explicit by incorporating the possibility of a sign switch into the metric. This is simply a notational convenience which allows them to keep the time coordinate real on both sides of the signature change surface. One could equally well leave the sign switch out of the metric (as in the unmodified Robertson-Walker form), in which case the change of signature would still take place when (according to Ellis, *et al.*'s proposed criterion) the cosmic dynamics caused da/dt to become imaginary. In this case, the signature change would manifest itself by a Wick rotation of the time coordinate from t to $\sqrt{-1}\tau$ rather than by a change of sign in g_{00} . It should be noted that in the physical (as opposed to coordinate artefact) signature change considered by Ellis, *et al.* and others, signature change occurs *either* by a change of sign of the metric *or* by a Wick rotation of the time coordinate, *but not both*. In Humphreys' coordinate-artefact-induced metric sign change, there is *both* a metric sign change *and* a Wick rotation of the time coordinate, and the two cancel each other, leaving the intrinsic signature of spacetime unchanged. The intrinsic signature change considered by Ellis, *et al.* is a coordinate-independent physical process which is caused by the dynamics of the cosmic expansion, while Humphreys' coordinate-artefact-induced sign change is not a physical process at all, but simply an artefact of the particular coordinate system he prefers to use, the Klein coordinate system.

- 35. Peebles, P.J.E., *Principles of Physical Cosmology*, Princeton University Press, pp. 312–313, 1993. This equation is strictly valid only for 'cold dust' cosmologies, but these are an excellent approximation to the actual universe throughout its observable ($Z < 1100$) history.
- 36. This transformation can also be derived by using the Ellis, *et al.* form of the metric to calculate the proper time interval elapsed on comoving clocks, in the same manner as is done for Humphreys' modified metric in note 37. This calculation shows that Ellis, *et al.*'s proposed criterion for classical metric signature change, imaginary expansion rate, is valid.
- 37. In fact, even Humphreys' proposed further generalization of the lapse function N to be a function of both t and η leads identically to the Robertson-Walker form of the metric. Humphreys' proposed generalization of the conventional Robertson-Walker metric is

$$ds^2 = c^2 N(\tau, \eta) d\tau^2 - a^2 \left[\frac{d\eta^2}{1 - k\eta^2} + \eta^2 (d\theta^2 + \sin^2 \theta d\phi^2) \right]$$

(equation 14, p. 201, reference 3.). This is actually no generalization at all, as the following analysis shows. If we consider the trajectory of a comoving clock, $d\eta = d\theta = d\phi = 0$, it follows that $ds^2_{\text{comoving}} = c^2 d\tau^2_{\text{comoving}} = c^2 N(\tau, \eta) d\tau^2$. This relation determines the mathematical form of the lapse function in terms of the comoving proper time interval dt_{comoving} and the coordinate time interval $d\tau$: $N(\tau, \eta) = d\tau_{\text{comoving}}/d\tau$. Substituting this formula for the lapse function $N(\tau, \eta)$ into Humphreys' modified metric immediately recovers the familiar Robertson-Walker form, which shows that the two equations are really the same.

- 38. Humphreys, D.R., *Starlight and Time: Solving the Puzzle of Distant Starlight in a Young Universe*, Master Books, Green Forest, Arkansas, pp. 10, 46, 1994, 1998.

Starlight and time: a response

D. Russell Humphreys

I thank Mr Conner and Dr Page for continuing to call attention to my little book on cosmology, *Starlight and Time*.¹ I often wonder if its persisting popularity is partly due to their determined attempts to discredit it. Of course, such a result would be far from what they desire, since their aim is to support Dr Hugh Ross's theistic evolutionary² version of the 'big bang' cosmology.

Another reason I am grateful for their critiques is that unsympathetic scrutiny, while not being particularly comfortable, either exposes flaws or, failing to do so, builds up confidence in the theory being scrutinized. I am happy to report that their latest attempt has had the latter effect, at least on me. That is because they have merely continued with Mr Conner's previous (1999) lines of attack,³ without paying adequate attention to my responses⁴ to those same arguments. Below I respond to these latest versions of their arguments, following the same order as in my 1999 reply.

They still have problems with a centre

In their 1998 critique,⁵ Conner and Page argued that both bounded-matter and unbounded-matter universes would have the same gravitational forces, so that there would be no essential difference between my cosmology and theirs. Their first step was to try to show that an infinite (unbounded) Newtonian cosmos uniformly filled with matter would have the same forces as a finite (bounded-matter) one. Here I have reproduced Figure 2(d) of their 1998 article, showing their result. The arrows show the pattern of gravitational forces they derived.

In my 1998 reply,⁶ I pointed out an alleged error in their derivation. In defence of their derivation, Conner and Page introduce several strained definitions. For example, they stretch out the meaning of the word 'centre' to include their idea of 'infinitely many' non-unique centres. But they seem to have missed my main point: a uniform unbounded-matter cosmos **cannot have a unique centre**. They seem to acknowledge this inadvertently by saying that the various D s in equations (1) through (3) are distances from the 'adopted' origin of coordinates. In Figure 2(d), they showed arrows of force converging upon a dot. The dot is the 'adopted origin of coordinates' caused by their method of analysis. Let's call it 'point C'. Here is the crucial problem with their result: **their 'forces' depend on where they choose to put point C**.

Point C is an arbitrary artefact of their method of analysis, existing only in the mind of the analyst. Another

analyst might place C in a different place. Yet the Newtonian cosmos they postulated is static, motionless on a large scale. That means the forces they derive should be measurable, and therefore physically real. For example, we could measure the directions of the forces with a plumb line. So how could the derived forces be physically real if they are to point toward a purely mental location? Would the plumb line change its direction if we were to change our mental placement of the ‘*adopted origin of coordinates*’? The answer is, no — something is clearly wrong with their derivation. Whether that flaw is the use of Newton’s ‘hollow shell’ theorem in a situation where it is not valid (as I alleged) is not the most relevant point. The most relevant point is that their **conclusion** is illogical. In writing my 1998 article, I had thought a quote from a cosmologist they respect would settle the issue:

*‘On the other hand, if matter were evenly dispersed through an **infinite** space, there would be **no center** to which it could fall.’⁶*

These are the words of Nobel Laureate Steven Weinberg⁷ (the second emphasis is mine). He was talking about precisely the situation Conner and Page were analyzing, an infinite uniform-matter Newtonian universe. Weinberg confirms what was already clear to me: without boundaries or variations in density, we can define **no unique centre** toward which gravitational forces could make matter fall. But Conner and Page do not agree to that point. In fact, nowhere does their letter take notice of Weinberg’s statement, though I quoted it in my 1998 paper and referred to it prominently in my 1999 reply. If Conner and Page cannot acknowledge such obvious features of a simple Newtonian theory, how can we have confidence in their pronouncements about much more subtle relativistic theories?

They still use circular reasoning about time

Referring to equation (20) of my 1998 article, the change in proper time τ measured by physical clocks at rest in the **centre** of a bounded-matter cosmos depends on the change in Schwarzschild (or coordinate) time t as follows:

$$d\tau^2 = \beta dt^2 \tag{1}$$

where β is Klein’s time dilation factor. In my 1998 article I pointed out that early in the expansion of the cosmos, β is negative, becoming positive only later. In the same article I offered reasoning that Schwarzschild time t is a *conceptual* coordinate, so that dt^2 remains positive throughout the expansion. Then the product βdt^2 would change sign during the expansion, being positive late in the expansion but negative early in the expansion. That would mean that $d\tau^2$ would be negative early in the expansion. The interval $d\tau$ would be imaginary (in the mathematical sense), or better, space like. As I pointed out, relativists interpret that as meaning that physical clocks would be stopped.

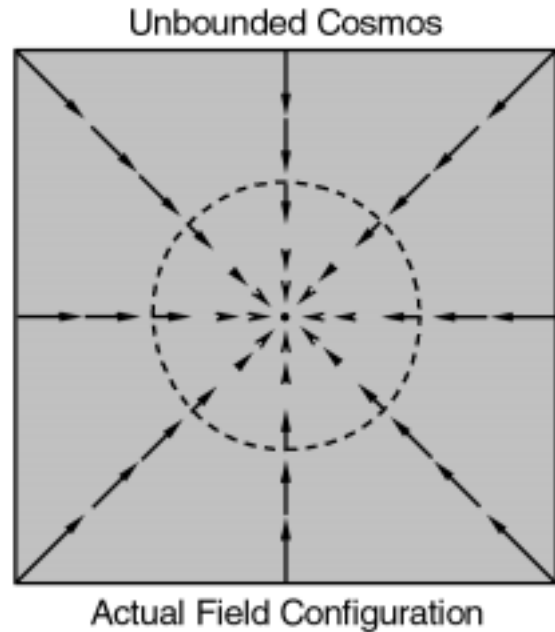


Figure 2(d) reproduced from 1998 critique by Conner and Page.⁵

My idea that the square of the Schwarzschild time interval dt^2 does not change sign is not new. It is the usual way relativists interpret the Schwarzschild metric in vacuum. Using my notation, that metric as applied to radial motions is

$$d\tau^2 = \beta dt^2 - \alpha dr^2 \tag{2}$$

In this metric, $\beta = 1 - (r_s / r)$, $\alpha = (c^2 \beta)^{-1}$, and r_s is the radius of the event horizon. Inside the event horizon both β and α are negative. Relativist theorists say this means that r for a particle inside the event horizon cannot be constant; otherwise dr^2 would be zero and the proper time interval $d\tau^2$ would become negative.⁸ Now if the square of the Schwarzschild time interval dt^2 could become negative, then $d\tau^2$ could be positive even if the particle’s r -coordinate were constant ($dr^2 = 0$). But relativists never seem to consider that option, perhaps because they interpret Schwarzschild time as a *conceptual coordinate*, at least unconsciously. This supports my reasoning that the square of the Schwarzschild time interval dt^2 remains positive through a signature change.

Conner and Page disagree. They are correct in saying that the time dilation factor β changes sign because (simplifying their notation a bit ⁹) the factor $(\partial\tau / \partial t)^2$ in their equation (7) changes sign. But they claim that it does so because it is the Schwarzschild (or ‘*Klein*’) time interval dt which becomes imaginary, not the proper time interval $d\tau$. However, they provide no proof here for their claim. They merely refer (their ref. 22) to a derivation in Conner’s 1999 letter.

My 1999 criticism of Conner’s derivation was that he

had based it on a questionable foundation. He assumed that his starting point, equation (1) in his 1999 letter, was valid in the timeless (Euclidean) zone. That is equivalent to assuming the point he wanted to prove; i.e. his reasoning was circular. But previously, in my 1998 article, I had questioned whether that very equation is valid in a Euclidean zone. In my 1999 reply, I emphasized those doubts. In their present letter, Conner and Page answer my criticism by using their conclusion to justify their starting point. **Circular reasoning** again! This leaves their case unproved.

My basic case for time dilation does not depend on the above point, the stopping of time in a Euclidean zone. As I asserted in my book and then pointed out on page 203 of my 1998 article, time dilation also occurs at the event horizon:

‘Therefore physical clocks at the centre of a white hole must stop (relative to Schwarzschild time) when the event horizon arrives.’

Contrary to an allegation by Hugh Ross,¹⁰ this quote shows I never gave up on that first possibility, time dilation at the event horizon.

Last year I came across a **new paper** which **supports** my view above. It was published in the *Astrophysical Journal* in 1995, only a year after my book was published.¹¹ The author, Martin Harwit, asserts that physical clocks near an event horizon tick slower than physical clocks which are far away from it. He refers not to Schwarzschild time, but to **proper** time in co-moving reference frames, the same sort of time and frames Conner and Page prefer. This means that their arguments about the meaning of Schwarzschild time are irrelevant to the question of time dilation at the event horizon.

They still misunderstand my model

Conner and Page’s reactions to some of the ‘signature change’ articles in secular relativity journals are useful to me; hitherto, most relativists have been fairly quiet about those developments. Here are my responses to the three comments in their sixth-from-last paragraph:

- (1) ‘... the speculative’ character of this literature’ ‘Speculative’ means different things to different people. For example, I think the currently popular string theories are highly speculative. But the basic observation by George Ellis — that Einstein’s field equations do not exclude the possibility of signature change — is on rock-solid ground. It is not at all speculative to try to explore the new territory Ellis has opened up. (Conner and Page’s parenthetical comment here merely repeats the conclusion of their circular reasoning, as I explained in section 2.)
- (2) ‘... criterion for signature change’ I have already been considering, and will continue to consider carefully, whether gravitational potential energy can

produce signature change. All writers have been rather unclear on precisely what would cause the changes, so I am not committed to any particular details of the picture I presented. We are indeed at the frontiers of human knowledge here, and I welcome knowledgeable instruction on these points.

- (3) ‘... applies to unbounded as well as bounded’ I certainly was not trying to say that unbounded universes could not have signature change. I merely was pointing out that bounded-matter universes have an additional factor to consider, namely gravitational potential energy.

The next three paragraphs, including equations (8) and (9), are merely a belaboured attempt to show that the Robertson-Walker metric can allow a signature change, even if one does not include an explicit lapse function. Okay, I’ll agree with that; I’m quite happy for them to now be allowing signature change. However if they had included a lapse function explicitly in their metric, they would have been more likely to see its effects in the equations. They did not do such.

In the second-to-last paragraph, Conner and Page assert that the cosmic microwave background radiation we see must have originated after any signature change. I agree. That is an implication of section 11 of my 1998 paper, and Figure 11 therein, in which the ‘light ray’ includes light from such sources.

In their last paragraph, they assert that a signature change would have to be simultaneous throughout the cosmos. However, their supporting sentence for this, ‘*This simultaneity is imposed by the fact that ...*’, turns out to rest on an assumption of the truth of the previous sentence — circular reasoning again! Their assertions about a 6,000 light-year particle horizon are built on the same inadequate logic. I exhort them to consider Figure 11 more carefully, especially on how it provides a counter example to their reasoning.

They still don’t acknowledge confirming research

Note carefully: in all their comments on the literature, Conner and Page have **completely ignored** my pointed challenge to comment on a 1997 paper by Hellaby *et al.*,¹² which asserted that such a timeless zone could occur in a black-hole/white-hole situation:

‘We have succeeded in demonstrating the possibility that a change in the signature of spacetime may occur in the late stages of black hole collapse, resulting in a Euclidean region which bounces and re-expands, passing through a second signature change to a new expanding Lorentzian space.’

Since that conclusion supports the main point of my 1998 paper, why do Conner and Page continue to remain silent about Hellaby’s 1997 paper?

Conclusions

In summary, Conner and Page have ignored the essence of all my 1999 challenges to them. Section 1 shows they are still not acknowledging their problems with a **unique centre** of the cosmos, heeding neither me nor Stephen Weinberg. Section 2 shows that they did not break out of their **circular reasoning** about the interpretation of the various time coordinates. Section 3 shows they are still attacking only strawman versions of my model. Section 4 shows they are still not acknowledging the most important **supporting paper**, the 1997 article by Charles Hellaby *et al.*

Their continued silence about the Weinberg and Hellaby quotes is very significant. Because Conner and Page have not contested my interpretation of those two quotes, the reader would be justified in considering their silence to be indirect support for my points. I will be interested to see how Conner and Page respond to the **new support** I cited from the literature, the 1995 *Astrophysical Journal* article by Harwit.

In all of this, let me emphasize that I am not claiming to be omniscient or inerrant! For example, I do not know whether my interpretation of Schwarzschild time is correct. I merely know that Conner and Page have not proved their case, and that they are ignoring the most important issues. Furthermore, I suspect there are mysteries related to the interpretation of time which no human yet understands. In general, I regard my work as one incomplete example of a new **class** of theories; centric cosmologies with various types of time dilation. I urge gifted creationists, who have the advantage of knowing from both Scripture and science that the world is young, to become expert in general relativity. I call upon them to generate better cosmologies than mine, to the glory of God.

References

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- By 'theistic evolutionism' I simply mean any attempt to reconcile **theism** with the events, sequence, and time scale alleged by naturalistic **evolutionism**, both physical and biological.
- Conner, S.R., Vistas — one more, *CEN Tech. J.* **13**(1):56–58, 1999.
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- Weinberg, S., *The First Three Minutes: A Modern View of the Origin of the Universe*, second paper edition, Basic Books (Harper-Collins), New York, p. 32, 1993.
- Rindler, W., *Essential Relativity: Special, General, and Cosmological*, revised second edition, Springer-Verlag, New York, p.149, 1977.
- Conner and Page change t to t_k and call it 'Klein' time, but here I am dropping their subscript to keep the notation simple. Their notation contains some psychological subtleties. Notice that the factor $(\partial t^{Klein} / \partial \tau)^2$ in their equation (7) is more complicated than necessary. They could have inverted the expression and dropped the cumbersome subscript, as I did in the text. Why didn't they clean up their notation? Probably because they want τ to be a dependent variable and t to be an independent variable. The usual mathematical convention is to subscript a dependent variable and put it in the numerator. It is okay to use notation to emphasize a point one wishes to prove, but here it would have been good to note that explicitly, since it is the very point at issue.
- Ross, H., Humphreys' new vistas of space, *CEN Tech. J.* **13**(1):49–50, 1999.
- Harwit, M., Time and its evolution in an inhomogeneous universe, *Astrophysical Journal*, **447**:482–490, 1995. Not being interested specifically in time at the event horizon, Harwit does not state its behaviour explicitly in words. However, his equation (11) relates (A) the rate of a clock moving with the surface of an expanding (or contracting) dust cloud to (B) the rate of a clock co-moving with the inner surface of an expanding shell of dust much further away. In the equation, when the radius r_{in} of the first clock becomes equal to the Schwarzschild radius r_s , the rate of the first clock becomes zero, as measured by the more distant second clock. This means that time dilation at the event horizon is a slowing of **proper** time as measured by **co-moving** physical clocks. It is not merely an alleged 'artefact' of using Schwarzschild time.
- Hellaby, C., Sumeruk, A. and Ellis, G.F.R., Classical signature change in the black hole topology, *Int. J. Modern Physics*, **D6**(2):211–238, 1997.